

## Lecture 9: Quantum walk search and collision finding

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## 1 Quantum walk search

Consider a graph  $G = (V, E)$  with  $|V| = n$  nodes and spectral gap  $\delta$ . Let  $M \subseteq V$  denote a subset of marked nodes of size  $|M| = m$ . Following the last exercise session, we have the following quantum walk search algorithm:

1. Set up the stationary state

$$|\pi\rangle = \frac{1}{\sqrt{n}} \sum_{x \in V} |\psi_x\rangle.$$

2. Repeat  $O(\sqrt{n/m})$  times:
  - (a) Reflect around marked subspace  $\text{span}_{x \in M}\{|\psi_x\rangle\}$  (i.e., apply  $2 \sum_{x \in M} |\psi_x\rangle \langle \psi_x| - I$ ).
  - (b) Reflect around stationary state  $|\pi\rangle$  (i.e., apply  $2 |\pi\rangle \langle \pi| - I$ ). We can do this using  $O(1/\sqrt{\delta})$  QW steps (see exercises).

The resulting state will have a constant overlap with the marked state  $|\pi_M\rangle = \frac{1}{\sqrt{m}} \sum_{x \in M} |\psi_x\rangle$ , so that measuring the state returns a marked element with constant probability.

We can summarize the (query) cost of the QW search algorithm using the following quantities:

- “Setup cost”  $\mathcal{S}$ : the number of queries needed to create the initial state  $|\pi\rangle$  in step 1.
- “Checking cost”  $\mathcal{C}$ : the number of queries needed to check whether an element is marked in step 2.(a).
- “Update cost”  $\mathcal{U}$ : the number of queries needed to implement a step of the QW.

The total cost (up to constants) is then

$$\mathcal{S} + \sqrt{\frac{n}{m}} \left( \mathcal{C} + \frac{1}{\sqrt{\delta}} \mathcal{U} \right).$$

This algorithm is called the “MNRS algorithm”, after Magniez-Nayak-Roland-Santha [MNRS07].

## 2 Collision finding with quantum walk search

Assume that we are given an array of integers  $x_1, x_2, \dots, x_N$ . A *collision* is a pair of distinct  $i, j$  such that  $x_i = x_j$ . How many elements do we have to query in order to find a collision (or decide that no collision exists)? Classically this essentially requires to query the full array, and so the classical query complexity is  $\Omega(N)$ . In contrast, using an algorithm based on *quantum walk search* we can find a collision with a *sublinear* number of queries.<sup>1</sup>

<sup>1</sup>While we focus on query complexity for ease of exposition, all algorithms can be implemented with a similar runtime.

The quantum walk algorithm for collision finding was proposed by Ambainis [Amb07]. The algorithm runs quantum walk search over elements or “words”  $\mathcal{Y} = (Y, x_Y)$ , consisting of (i) a size- $k$  subset  $Y \subseteq [N]$ , and (ii) the list  $x_Y$  of integers  $x_j$  with index  $j \in Y$ . We call an element  $\mathcal{Y}$  marked if  $Y$  contains both indices of a collision (equivalently,  $x_Y$  must contain a collision).

**Exercise 1.** Let  $n = \binom{N}{k}$  denote the number of elements and  $m$  the number of marked elements. Show that  $m/n \in \Omega(k^2/N^2)$ .

To use quantum walk search, we consider a graph  $G$  with vertex set  $V$  indexed by the elements  $\mathcal{Y}$ . There is an edge between  $\mathcal{Y} = (Y, x_Y)$  and  $\mathcal{Y}' = (Y', x_{Y'})$  if the subsets  $Y$  and  $Y'$  differ in exactly one element (i.e., we can obtain  $Y'$  from  $Y$  by replacing one index). The resulting graph  $G$  has  $n = \binom{N}{k}$  vertices and is  $k(n-k)$ -regular. It corresponds to a so-called *Johnson graph*, and one can show that its spectral gap is  $\delta \in \Omega(1/k)$  when  $k \ll n$ .

A star state  $|\psi_{\mathcal{Y}}\rangle$  centered on a vertex  $\mathcal{Y}$  of  $G$  is given by the state

$$|\psi_{\mathcal{Y}}\rangle = \frac{1}{\sqrt{k(n-k)}} \sum_{\mathcal{Y}' \sim \mathcal{Y}} |\mathcal{Y}, \mathcal{Y}'\rangle,$$

where the sum runs over neighboring elements  $\mathcal{Y}'$  of  $\mathcal{Y}$ . The quantum walk search algorithm then starts from the uniform superposition

$$|\pi\rangle = \frac{1}{\sqrt{n}} \sum_{\mathcal{Y}} |\psi_{\mathcal{Y}}\rangle,$$

and the algorithm has cost

$$\mathcal{S} + \frac{N}{k}(\sqrt{k}\mathcal{U} + \mathcal{C}),$$

where  $\mathcal{U}$  is the *update cost* or cost of implementing a single quantum walk step on the Johnson graph  $G$ . We now bound the different costs.

For the checking cost  $\mathcal{C}$ , note that we can check whether a given state  $|\psi_{\mathcal{Y}}\rangle$  is marked simply by checking whether the list  $x_Y$  contains a collision. Since this list is given explicitly in the description of  $|\psi_{\mathcal{Y}}\rangle$ , this requires no queries and so  $\mathcal{C} = 0$ .

The setup cost  $\mathcal{S}$  amounts to creating the state  $|\pi\rangle = \frac{1}{\sqrt{n}} \sum_{\mathcal{Y}} |\psi_{\mathcal{Y}}\rangle$ . We do this in a few steps. First, we prepare the state  $\frac{1}{\sqrt{n}} \sum_{\mathcal{Y}} |\mathcal{Y}\rangle |0\rangle$ . This takes  $k$  queries. Then, we construct the mapping  $U_{\psi}$  defined by  $U_{\psi} |\mathcal{Y}\rangle |0\rangle = |\psi_{\mathcal{Y}}\rangle$ . We do this in two steps:

$$\begin{aligned} |\mathcal{Y}\rangle |0\rangle &= |Y, x_Y\rangle |0\rangle \stackrel{(i)}{\mapsto} \frac{1}{\sqrt{k(n-k)}} \sum_{\mathcal{Y}' \sim \mathcal{Y}} |Y, x_Y\rangle |Y', 0\rangle \\ &\stackrel{(ii)}{\mapsto} \frac{1}{\sqrt{k(n-k)}} \sum_{\mathcal{Y}' \sim \mathcal{Y}} |Y, x_Y\rangle |Y', x_{Y'}\rangle = |\psi_{\mathcal{Y}}\rangle. \end{aligned}$$

Step (i) requires no queries. Step (ii) amounts to gathering the elements  $x_{Y'}$  with index in  $Y'$ . Since  $x_{Y'}$  contains exactly one element not in  $x_Y$ , this requires exactly one query. The setup cost  $\mathcal{S}$  is hence roughly  $k$ .

Finally, we bound the cost  $\mathcal{U}$  of a single call to the quantum walk operator  $W$ . Recall from last lecture that  $W = S \cdot C$  where  $S$  is a simple swap (i.e.,  $S |\mathcal{Y}, \mathcal{Y}'\rangle = |\mathcal{Y}', \mathcal{Y}\rangle$ ), requiring no queries, and  $C = 2(\sum_{\mathcal{Y}} |\psi_{\mathcal{Y}}\rangle \langle \psi_{\mathcal{Y}}|) - I$  is a reflection around the star subspace. Using a similar trick as before, we can implement the reflection  $C$  by making two calls to the preparation operator  $U_{\psi}$ , which requires a single query. This proves that the update cost  $\mathcal{U}$  is  $O(1)$ .

**Exercise 2.** Verify that  $C = U_\psi R_0 U_\psi^\dagger$ .

Combining these different arguments, we can bound the total cost by

$$\mathcal{S} + \sqrt{\frac{n}{m}} \left( \frac{1}{\sqrt{\delta}} \mathcal{U} + \mathcal{C} \right) \approx k + \frac{N}{\sqrt{k}}.$$

If we set  $k = N^{2/3}$  then this yields a quantum algorithm for collision finding with complexity  $\tilde{O}(N^{2/3})$ . This is essentially optimal by the  $\Omega(N^{2/3})$  lower bound of Aaronson and Shi [AS04].

## References

- [Amb07] Andris Ambainis. Quantum walk algorithm for element distinctness. *SIAM Journal on Computing*, 37(1):210–239, 2007.
- [AS04] Scott Aaronson and Yaoyun Shi. Quantum lower bounds for the collision and the element distinctness problems. *Journal of the ACM (JACM)*, 51(4):595–605, 2004.
- [MNRS07] Frédéric Magniez, Ashwin Nayak, Jérémie Roland, and Miklos Santha. Search via quantum walk. In *Proceedings of the thirty-ninth annual ACM symposium on Theory of computing*, pages 575–584, 2007.